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40 3 RADIOFREQUENCY ELECTROMAGNETIC FIELDS INSIDE THE BODY

41 The body interacts with radiofrequency (RF) electromagnetic fields (EMFs) and the strengths of the
42 fields inside the body are quite different from those outside.

43 This chapter describes the relationships between external fields to which the body is exposed and the
44 induced fields that result inside the body. The scope is mainly restricted to studies in humans, although the
45 methods and principles described apply equally to characterising the exposures of animals, cells and other
46 biological structures in experimental systems.

47 3.1 Principles of interactions with the body

48 A radio frequency (RF) electromagnetic field (EMF) in air can be reflected, transmitted, refracted or
49 scattered by a biological body. The reflected and scattered fields may proceed in directions different from that of
50 the incident RF field, while the transmitted and refracted fields interact (are coupled) with the biological body in
51 different ways. These interactions are strongly dependent on the frequency, waveform, and strength of the
52 induced fields as well as the energy deposited or absorbed in the biological system. In addition, the distribution
53 of the fields inside a biological system such as the human body is affected by the distance and location of the
54 source with respect to the body, the anatomy, posture and the surrounding environment of the body.

55 3.1.1 Coupling of the body to the field

56 Coupling between an incident uniform (plane-wave) RF field and the human body depends only on
57 the frequency, the direction of incidence and the polarisation of the field. Where exposure to non-uniform fields
58 occurs, as when in close proximity to sources, the coupling also depends on the source characteristics, position
59 and orientation with respect to the body. Source-related considerations affecting exposure, such as near and far-
60 field regions, and mutual coupling, are discussed in Chapter 2, as are environmental considerations such as
61 reflections and multipath propagation.

62 Coupling is strongly dependent on the ratio of wavelength, λ , to that of the body dimensions; therefore
63 it is convenient to divide the RF EMF spectrum into three regions when considering physical interactions.

- 64 1. Body dimensions are small in relation to the wavelength, λ , for frequencies between 100 kHz and
65 10 MHz ($\lambda > 30$ m). Coupling to the body is weak, but RF energy penetrates deeply so generalised
66 absorption occurs throughout the body tissues.
- 67 2. Body dimensions such as height and limb length become comparable with the wavelength for
68 frequencies in the range 10 MHz to 2 GHz ($30 \text{ m} > \lambda > 15 \text{ cm}$) so resonances can occur, giving a
69 strong dependence between coupling and frequency and a complicated pattern of absorption in the
70 body tissues.
- 71 3. Between 2 GHz and 300 GHz ($\lambda < 15 \text{ cm}$) the wavelength is small in relation to the body tissues and
72 coupling is characterised by small penetration into the body tissues and surface dominated absorption.

73 These aspects of coupling and the mechanisms giving rise to them are discussed further in Section 3.1.4.

74 Both direct and indirect coupling of external RF fields into the human body can occur (see next
75 section) and result in the induction of fields and currents inside the body tissues. Dielectric and ohmic losses in
76 the tissues lead to energy absorption and a rise in local or whole body temperature. In addition, where the
77 exposure occurs within the reactive near-field region with respect to a source (see chapter 2) the electric and
78 magnetic components of the incident field interact differently with the body. Hence, both quantities and their
79 respective interactions must be determined separately and summed together in order to fully characterise the
80 human exposure.

81 3.1.2 Direct and indirect effects

82 RF EMFs carry energy and when the body is exposed to them, some of the energy is absorbed, a
83 *direct effect* which leads to heating of the body tissues (WHO, 1993). At frequencies below 100 kHz (and
84 therefore not in the scope of this document), the physical quantity identifiable with most biological effects is the
85 electric field strength in tissue, which is related to the current density. These effects are related to electrical

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86 stimulation of tissues, but become less prominent as frequency increases and are disregarded for frequencies
87 above 10 MHz. Generally, in the RF region (and exclusively above 10 MHz), the rate at which the body is
88 heated is considered a more appropriate measure to assess the exposure.

89 Contact with a radio antenna or metallic conductor placed in a RF field can lead in some
90 circumstances to electric shock or burn, *indirect effects* that result from current flow in the body tissues
91 (Chatterjee, Wu & Gandhi, 1986). Shock is related to electrostimulation of tissues and burn occurs due to intense
92 and rapid localised heating. Electrical burns follow the path of current flow through the body tissues and can be
93 much deeper than burns that result from contact with hot objects. They can occur at points where current exits
94 the body as well as where it enters the body.

95 Many factors determine the potential for electrical burn or shock including the power density,
96 frequency of the signals, the grounding conditions, whether the structure has resonant dimensions and how much
97 of the body is in contact with the conductor. Electrical burns can be a hazard for those working near antennas
98 and other metallic structures either radiating or exposed to RF fields. However, they are exploited to beneficial
99 effect in medicine where they are applied in electrosurgery.

100 **3.1.3 Absorption of RF energy**

101 The absorption of energy from RF EMF causes molecules to vibrate which in turn leads to heating of
102 body tissues. This absorption (and the consequent heating) is defined by a quantity known as the specific energy
103 absorption rate (SAR), with unit, watt per kilogram (W/kg). SAR is derived from the square of the
104 (instantaneous) electric field strength, E , in tissue:

$$105 \quad \text{SAR} = \frac{1}{2} \frac{\sigma}{\rho} E^2 \quad (3.1)$$

106 where σ and ρ are the conductivity (in siemens per metre, S/m) and the density (kg/m^3) of the tissue of interest.
107 For a sinusoidally varying electric field, the factor of $\frac{1}{2}$ may be omitted and the rms (root mean squared, see
108 Chapter 2) value of the field substituted in the above equation in order to obtain the time-averaged SAR.

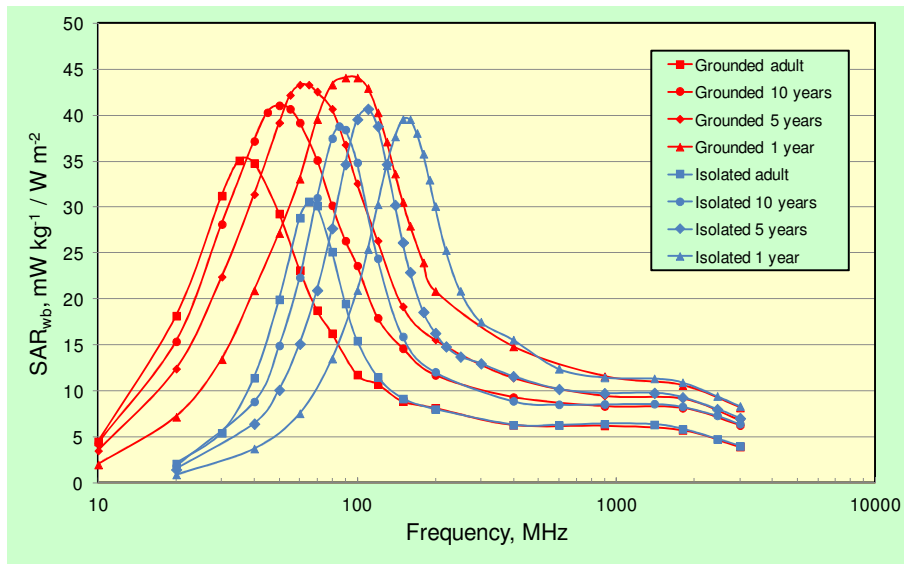
109 SAR provides a measure of the power absorbed from a radio frequency signal per kilogram of body
110 tissue and is often used as a proxy for the amount of heating or temperature rise in the body. It may be derived at
111 a point in time and space in the body tissues; however, it is more usual to average the quantity over time and
112 space in some appropriate way. Whole-body SAR may be derived by averaging the quantity over all tissues of
113 the body or localised SAR may be derived for a particular organ, tissue type or body part, e.g. by averaging over
114 a 1 g or 10 g mass of tissue in which the maximum SAR occurs. Averaging time is relevant to thermal
115 considerations because the heating potential of an exposure depends on how quickly a certain amount of energy
116 can be input in relation to how quickly it can be dissipated by the body. SAR is generally averaged over a period
117 of 6 minutes in order to derive a quantity closely related to the heating potential of absorbed RF energy
118 (ICNIRP, 1998).

119 SAR cannot be measured non-invasively so it cannot be measured easily in a living system. It is
120 therefore usually estimated from simulations using experimental or computer-based models of the living body
121 (see section 3.2).

122 **3.1.4 Factors affecting human exposure**

123 **3.1.4.1 Body resonance**

124 The human body acts as an oblong mass of poorly conducting material when exposed to RF fields. It
125 is particularly effective at absorbing RF fields under resonant conditions where the wavelength of the radiation is
126 comparable with the dimensions of the body. Maximum coupling occurs if the incident electric field is polarized
127 parallel to the long body axis and resonant currents flow up and down the body. The resonance occurs when the
128 body height is approximately half a wavelength ($\lambda/2$) for an ungrounded body and nearer a quarter wavelength
129 ($\lambda/4$) for a person standing in electrical contact with a conducting ground plane (Conil et al., 2008; Dimbylow,
130 2005a; Dimbylow, 1997; Kühn et al., 2009). An example set of computational results showing how SAR
131 averaged over the whole body mass varies in relation to body height and grounding conditions is shown in
132 Figure 3.1.



133

134 **Figure 3.1.** Predicted whole body SAR in typical models of the human body per unit power density of a
 135 vertically polarised plane wave incident towards the chest. The basic model is of a 73 kg, 1.76 m tall adult male
 136 standing with its arms to the side and this has been scaled to have heights/masses representative of children at
 137 different ages. Conditions are also considered where the shoes effectively form an electrical contact between the
 138 body and a conducting plane beneath it and where they provide isolation (Data from Dimbylow (2005a; 1997).

139 The frequency of resonance depends on the size, shape and posture of the exposed individual. It also
 140 depends on the direction from which the wave is incident and the polarisation of the incident electromagnetic
 141 field with respect to the body. In the case of a vertically polarised wave incident on the body, as shown in Figure
 142 3.1, the resonance frequency is particularly strongly affected by the height of the exposed person and the ease
 143 with which current can flow through the feet to ground. For standing adults, the peak of this resonant absorption
 144 occurs in the frequency range of 70–80 MHz if they are electrically isolated from ground and at about half this
 145 frequency if they are electrically grounded. Smaller adults and children show the resonance at higher
 146 frequencies. In addition to whole-body resonance, it is possible for partial-body resonances to occur, e.g. in the
 147 limbs.

148 **3.1.4.2 Exposure below body resonance**

149 For uniform (plane wave) exposures, as frequency reduces below that of the body resonance the body
 150 acts as a short poor conductor and couples increasingly weakly to RF fields, as shown in Figure 3.1. In this
 151 region, the SAR is approximately proportional to the height of a person and to the frequency squared.

152 In this frequency range, the highest exposures tend to occur in practice from near-field sources (see
 153 Chapter 2) that generally have strong field gradients and are used within a few metres of the body. Under such
 154 conditions, the energy is either capacitively or inductively coupled to the body (instead of through radiation),
 155 depending on whether electric or magnetic field component is the dominant source. Examples of such exposure
 156 scenarios are with dielectric heaters and diathermy applicators where electric field source is dominant, and with
 157 inductive cooking hobs, anti-theft systems, wireless power transfer systems, and magnetic resonance imaging
 158 where the magnetic field source is dominant.

159 **3.1.4.3 Exposures above body resonance**

160 At frequencies above body resonance, the body acts as a dielectric¹ object that is large with respect to
 161 the wavelength and to the penetration depth (skin depth – see below) of the radiation into the body tissues.
 162 Therefore, the absorbed energy is approximately proportional to the exposed surface area of the body (Gosselin
 163 et al., 2009). In turn, whole body SAR is proportional to the ratio of exposed surface area to the body mass

¹ Dielectric materials are materials in which locally bound electric charges can be displaced in response to an externally imposed field, leading to polarisation of domains inside the material.

164 (Kühn et al., 2009). This is consistent with the results shown in Figure 3.1, which show that children have higher
 165 whole-body SAR than adults at frequencies around and above body resonance.

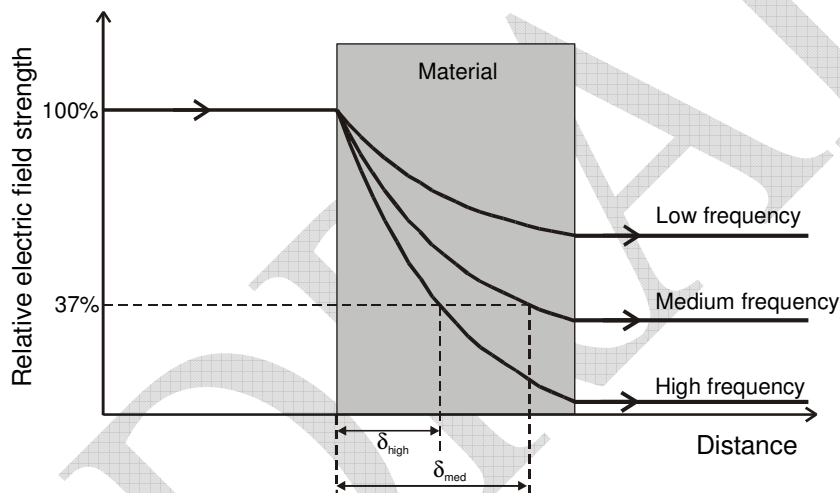
166 3.1.4.4 Skin depth

167 RF fields become less penetrating into body tissues as frequency increases so the energy they deposit
 168 in the body becomes increasingly confined to the body surface. A useful figure of merit to describe the
 169 penetration is the skin depth, a distance over which the electromagnetic field component of a wave penetrating
 170 into a material reduces to 37% of its initial value, as illustrated in Figure 3.2. The skin depth of tissues depends
 171 on their electrical permittivity² and conductivity. The general expression for skin depth for poor conductors
 172 (non-metals) at radio frequencies is as follows (Griffiths, 1989):

$$173 \quad \delta = \left(\frac{1}{\omega}\right) \left\{ \left(\frac{\mu\epsilon}{2}\right) \left[\left(1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2\right)^{1/2} - 1 \right] \right\}^{-1/2} \quad (3.2)$$

174 where ω is the angular frequency, ϵ , σ and μ are the permittivity (F/m), conductivity (S/m), and magnetic
 175 permeability of the materials respectively. In biological materials, μ in tissues has essentially the same value as
 176 that of free space, $4\pi \times 10^{-7}$ H/m.

177 The skin depths of tissues with low water content such as fat and bone are greater than those with
 178 higher water content such as muscle and skin. Table 3.1 contains typical skin depths for low and high water
 179 content tissues at selected frequencies.



180
 181 **Figure 3.2.** Illustration of how RF energy is absorbed in biological materials showing how the skin depth
 182 decreases with increasing frequency. Reflection of the incident radiation is assumed negligible at each interface
 183 in this diagram. The skin depth at high frequency, δ_{hi} , is less than that for medium frequency, δ_{med} .

184 Skin depth is also strongly frequency dependent. As is clear from Table 3.1 at around 150 MHz, bone
 185 and fat are poor conductors so they absorb energy weakly and the RF penetrates deeply into these tissues. High
 186 water content tissues such as muscle and skin on the other hand are good conductors and these will absorb more
 187 strongly and thus have a lower skin depth. At higher frequencies, the skin depth decreases, therefore absorption
 188 in the body becomes increasingly confined to surface tissues.

189 At mobile phone frequencies, of the order of 1 GHz or so, the skin depth in the brain is a few cm.
 190 Consequently most of the energy from the incident radiation is absorbed in one side of the head within a few
 191 centimetres of the handset. At 10 GHz the skin depth in most tissues is a few millimetres so almost all the energy

² For biological tissues in the radiofrequency range, this is generally a complex quantity, i.e. it has real and imaginary components that must both be taken into account.

192 will be absorbed in the skin and other surface tissues and there will be very little penetration into the deeper
 193 tissues of the body.

194

Table 3.1. The conductivity and skin depth of low and high water content tissues at selected radio frequencies*

Frequency	Tissues with low water content				Tissues with high water content			
	Fat		Bone		Muscle		Skin	
	Conductivity (S/m)	Skin depth (mm)	Conductivity (S/m)	Skin depth (mm)	Conductivity (S/m)	Skin depth (mm)	Conductivity (S/m)	Skin depth* (mm)
150 MHz	0.04	366.1	0.07	301.0	0.7	67.2	0.5	85.0
450 MHz	0.04	301.9	0.10	202.2	0.8	51.3	0.7	52.9
835 MHz	0.05	252.0	0.14	139.5	0.9	43.5	0.8	41.5
1.8 GHz	0.08	157.1	0.28	66.7	1.3	29.2	1.2	28.3
2.45 GHz	0.10	117.1	0.39	45.8	1.7	22.3	1.5	22.6
3 GHz	0.13	93.6	0.51	35.2	2.1	18.0	1.7	18.9
5 GHz	0.24	49.4	0.96	17.7	4.0	9.3	3.1	10.5
10 GHz	0.58	19.6	2.13	7.3	10.6	3.3	8.01	3.8

*The skin depth data in are calculated based on permittivity and conductivity of tissues taken from Gabriel et al. (1996b). The formula used for calculation of skin depth is taken from Griffiths (1989).

195

196 3.1.4.5 Type of tissue

197 The extent of reflection, absorption and transmission from an incident RF field depends on the type of
 198 material and its thickness in relation to the wavelength. When RF fields are incident on highly conducting (e.g.
 199 metal) surfaces, reflection is the dominant process; some absorption will occur and there will be almost no
 200 transmission. In the case of biological materials, there will be some reflection and the relative proportions of
 201 absorption and transmission will depend on the thickness of the material. Radio waves at telecommunications
 202 frequencies generally penetrate into the body tissues for a few centimetres before having been almost completely
 203 absorbed by the tissues (Table 3.1).

204 Absorption of energy also depends on the electrical properties of tissues which vary, mainly because
 205 some tissues are better electrical conductors than others. On exposure to RF fields, energy is not deposited
 206 uniformly throughout the body, even if the incident radiation has uniform power density. When radio waves are
 207 incident on a homogeneous slab of material, the proportion of energy transmitted decreases exponentially with
 208 the thickness of the slab. Consequently more energy will be absorbed in the material towards the front surface
 209 facing the incoming waves than towards the rear.

210 3.1.4.6 Heterogeneity of tissues

211 Biological systems such as the human body contain many different types of tissues that connect and
 212 interface with each other in many different conformations. Some tissues are extensively folded around each
 213 other, e.g. in the case of blood vessels within brain tissue, while others connect at smooth boundaries over large
 214 areas, e.g. the surfaces of bones. With such complex geometries, it is possible for internal focussing and
 215 reflection of RF fields to occur. Internal resonances can also occur where structures have comparable size to the
 216 wavelength. These processes lead to a very inhomogeneous distribution of absorbed RF energy even when the
 217 external fields to which the body is exposed are uniform. Only anatomically realistic computational models of
 218 the body are able to account for these physical processes and yield accurate predictions of absorbed energy at a
 219 detailed anatomical level. Such models are described later in this chapter.

220 3.2 Dosimetry

221 As explained above, the strength of the induced electric and magnetic fields inside the body are
 222 different from those outside. To assess the exposure of the body to an external EMF, one needs to determine

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223 internal quantities (also called dosimetric quantities) which are related to the exposure in tissue (e.g. the induced
224 electric field strength, induced current density and SAR). The process of determining internal quantities
225 (dosimetric quantities) from external fields (exposure quantities) is called dosimetry. The role of dosimetry is to
226 evaluate the induced quantities in the body and to correlate them with the biological effect of concern.

227 Chapter 4 of this document contains detailed discussions on different interaction mechanisms between
228 RF EMF and the human body. The most recognised and well established mechanism for biological effects of RF
229 radiation is tissue heating (Sheppard, Swicord & Balzano, 2008). Therefore, the current international guidelines
230 for human exposure (ICNIRP, 1998; IEEE, 2005) continue to be based on classical heating mechanisms and the
231 avoidance of adverse effects that occur in relation to temperature rises. The dosimetric quantity of most
232 relevance at frequencies up to 10 GHz, i.e. where RF energy penetrates appreciably into the body tissues, is
233 therefore the SAR, which describes the rate at which energy is absorbed in tissues per unit mass. At frequencies
234 above 10 GHz, incident power density is the dosimetric quantity that is most relevant because penetration of RF
235 energy into the body tissues is very small. As frequency reduces below 10 MHz, nerve and muscle stimulation
236 effects due to induced currents in the body become increasingly relevant. Therefore current density and/or
237 induced electric field strength are also used as dosimetric quantities at frequencies below 10 MHz (ICNIRP,
238 1998).

239 Presently used guidelines and standards recommend a set of *basic restriction* values in terms of SAR
240 (ICNIRP, 1998; IEEE, 2005). However, since it is not possible to measure SAR directly, the guidelines also
241 define *reference levels* that are used in practical assessment of compliance with the basic restrictions under a real
242 exposure scenario. These have been conservatively derived from the basic restrictions with the objective that if
243 the reference levels are not exceeded, the basic restrictions are also observed. Dosimetry plays an important role
244 in the implementation of guidelines especially in the derivation of the reference levels. Dosimetry is also
245 important when the exposure exceeds the reference levels. In this case it is necessary to examine whether the
246 exposure actually exceeds the basic restriction or not by means of dosimetry.

247 Dosimetry plays an important role in scientific research to assess the exposure of people to RF EMF.
248 The International EMF project of WHO emphasized the importance of well-defined exposure conditions for
249 biological experiments in order to achieve meaningful interpretation and reproducibility of the result (Repacholi,
250 1998). Kuster and Schönborn (2000) published the minimum requirements of exposure systems for biological
251 experiments addressing health effects of RF exposure.

252 There are different dosimetric methods available; experimental and numerical (theoretical) dosimetry
253 techniques can be used to assess internal fields for different sources and geometries. Each method has
254 advantages and disadvantages. For example, numerical dosimetry using realistic biological models can provide
255 very fine spatial distribution of SAR and induced current density. However, actual exposure conditions from a
256 real source can only be assumed in numerical dosimetry because the source has to be modelled. On the other
257 hand, there are practical limits to the anatomical detail that can be incorporated into physical models and tissue
258 boundaries would be disturbed if a measurement probe were manipulated inside them to measure the induced
259 fields. Both methods are described in detail in the following sections. Usually, scientists, select one of the
260 dosimetry techniques suitable for a particular exposure scenario, and validate the evaluated dose quantities by
261 comparing results between numerical and experimental dosimetry.

262 **3.2.1 Assessment of SAR**

263 The SAR inside the body exposed to RF fields can be assessed by dosimetric tools; namely
264 experimental and numerical techniques. Depending on the exposure scenario and the position of the RF field
265 source with respect to the body, whole body SAR or localised SAR values can be obtained and compared with
266 basic restriction values in the guidelines.

267 Significant temperature rise could occur if body parts are exposed to extremely localised exposures,
268 leading in thermal injury of the tissue irrespective of the deep-body temperature elevation. In these situations, the
269 local SAR in the part of the body is considered as the dosimetric quantity of interest. Temperature elevation in
270 different body parts, however, is not necessarily proportional to the local SAR because of the heat conduction
271 and transportation of heat by blood flow. Therefore, estimation of local temperature rise as an additional
272 dosimetric quantity can also sometimes also be undertaken.

273 **3.2.2 Mass-averaged SAR**

274 Rise in deep (core) body temperature is related to the “average” energy absorption rate in the body,
275 whereas temperature rise in organs, like the eyes and brain, depends primarily on the energy absorption in these
276 organs (ICNIRP, 1998). Therefore, SAR is usually averaged either over the whole body mass, or over a small
277 sample volume or mass of tissue, depending on the spatial pattern of the energy absorption in the body and
278 where temperature rises are expected.

279 The averaging mass serves two purposes: firstly, it provides a definition of SAR that is robust and not
280 overly sensitive to slight changes in exposure set-up; secondly, it takes into account the spatial dispersion of the
281 deposited energy which is provided by heat transfer mechanisms (NRPB, 2004a). The SAR values are highly
282 dependent on the geometry of the part of the body exposed and on the exact location and geometry of the source.

283 For frequencies up to a few GHz, as used in wireless communications, SAR is normally averaged over
284 either 1 or 10 g. The averaging mass should be large enough to maximise the correlation with local temperature
285 elevation (Hirata, Shirai & Fujiwara, 2008). It is suggested that an averaging mass of 10 g is a more appropriate
286 parameter when the correlation between SAR averaging mass and the elevated temperature is investigated
287 (Hirata, Shirai & Fujiwara, 2008; Razmadze et al., 2009).

288 The dominant factors influencing the correlation between mass-averaged SAR and temperature
289 elevation are the thermal diffusion length³ in the biological tissue, which largely depends on the blood perfusion
290 rate, and the penetration depth of the RF waves (Hirata & Fujiwara, 2009).

291 One of the rationales for the 10 g averaging mass is based on the temperature elevation in the brain
292 and eye lens, which has been computed with anatomically-based head models. The choice of averaging mass is
293 clearly most important for highly localised exposure. At the other extreme, a uniform SAR distribution gives the
294 same spatial average whatever mass is used. The greatest differences between 1 g and 10 g values are expected
295 for situations of near-field exposure to a small antenna such as that of a mobile phone (NRPB, 2004a). Also,
296 these averaging masses may not be appropriate for estimating temperature rises close to metallic medical
297 implants where significantly higher SAR values are calculated when it is averaged over 1 g (Kyriakou et al.,
298 2012; McIntosh, Iskra & Anderson, 2014).

299 **3.3 Biological models**

300 **3.3.1 Physical and numerical phantoms**

301 For health risk assessment, it is necessary to directly evaluate SAR, possibly from temperature rise, or
302 to evaluate SAR from the induced electric field strength in a human body exposed to RF EMFs. However, it is
303 very difficult to measure the internal electric field strength or temperature elevation in the human body using
304 non-invasive methods. Therefore a surrogate of the human body, a so called “phantom” is used. A phantom
305 designed to replicate the human body should exhibit similar interaction mechanisms with RF to those of a real
306 human body. It can be a physical model, made from different organic and non-organic materials to replicate the
307 electrical properties of various human tissues. It can also be a numerical model, which is an anatomically
308 realistic computer model of typical people or animals, with values of the electrical properties for the different
309 simulated tissues assigned to them. However, such phantoms generally do not include blood flow, which is an
310 important heat transportation mechanism in living systems. For this reason, temperature rises derived in such
311 phantoms will generally be higher than those produced in the body.

312 Physical phantoms are used in experimental dosimetry; the set-up usually consists of other
313 components such as emulators that represent the base stations to which sources used near the body would
314 normally connect with, implantable electric field sensors, a spatial scanning system, remote control and data
315 recording systems. Alternatively, the distribution of induced electric field strength and the corresponding SAR
316 values in the body can be estimated by solving Maxwell’s equations utilising different mathematical methods
317 such as finite difference time domain (FDTD) calculations (Tavlove & Hagness, 2005).

³ Thermal diffusion length is the distance range from the point of maximum temperature elevation within which the majority of thermal energy transfer takes place. This is comparable to the concept of skin depth when electromagnetic waves are absorbed in a conductive medium, and is the distance at which the temperature elevation has reduced to 1/e (i.e. 63%) of its value at the point of maximum elevation.

318 **3.3.2 Models of animal and human body**

319 *3.3.2.1 Experimental models*

320 Phantoms for RF dosimetry are required to simulate the electrical properties equivalent to those of the
321 human body. Various types of materials have been developed for phantoms and their references may be found in
322 international standards on RF dosimetry (IEC, 2005; 2002).

323 Phantoms made from organic materials have been widely used for RF dosimetry. Water and
324 gelatine are amongst the main ingredients in phantom recipes as they are relatively easy to prepare and to adjust
325 for their electrical properties. They are also suitable for SAR measurements where an electric field probe is
326 inserted into the phantom for scanning. On the other hand, these materials have poor stability of their electrical
327 properties due to water evaporation. Attempts have been made to make dry phantoms with high stability;
328 however they require complex and skilled procedures and are costly (Kobayashi et al., 1993; Nikawa, Chino &
329 Kikuchi, 1996).

330 Electrical properties are frequency-dependent and also vary according to the tissue's composition and
331 structure such that homogenized tissues do not necessarily have the same electrical properties as those with
332 cellular structures etc preserved intact. High-water-content tissues such as muscle are usually made with wet
333 material whereas recipes for low-water-content tissues such as fat and bone are made with dry material.

334 Each phantom recipe can simulate the electrical properties of a tissue from several hundred MHz to
335 several GHz (Hartsgrove, Kraszewski & Surowiec, 1987; Okano et al., 2000). It is however difficult to adjust the
336 electrical properties of the phantom within small deviation, e.g. 5 %, from those of the actual biological tissues
337 over broad frequency ranges. Different recipes optimized to the target electrical properties at each frequency are
338 therefore used for critical measurements such as SAR compliance tests (IEC, 2005).

339 Several studies have attempted to develop broad-band phantoms (Gabriel, 2007; Lazebnik et al., 2005;
340 Youngs et al., 2002). However, in adjusting both real part and imaginary part of the complex permittivity of the
341 phantom to the target values simultaneously, the uncertainty of the electrical properties measured by
342 commercially available systems, temperature change and water evaporation are amongst the difficulties that need
343 detailed consideration (Fukunaga, Watanabe & Yamanaka, 2004).

344 It is also very difficult to develop a heterogeneous structure using organic materials such as water and
345 gelatine. Moreover, the boundaries in such structures would be disturbed if a measurement probe were
346 manipulated through them or in their vicinity. Therefore, homogeneous tissue is generally used as a liquid inside
347 shaped hollow shells in order to form physical phantoms. For instance, a standardised phantom in the shape of a
348 hollow head is usually used for compliance tests of cellular phones (IEC, 2005; IEEE, 2003). A full-size model
349 of the human body was first developed by Olsen (1979) and by Olsen and Giner (1989). Stuchly et al. (1987)
350 developed a whole-body phantom which simulates heterogeneous structure with solid material for bone within a
351 liquid phantom for high-water content tissues such as muscle. They measured electric field distributions by
352 scanning with a field probe inside the heterogeneous phantom. Several heterogeneous head phantoms have also
353 been developed for SAR evaluation during use of a mobile wireless handset (Cleveland & Athey, 1989; Okano
354 et al., 2000).

355 *3.3.2.2 Numerical models*

356 Simple models such as spheres and ellipsoids are amenable to analytical solutions and can be used to
357 characterize the RF energy absorption in a human body. Durney et al. (1986) systematically reviewed and
358 summarized these models and their use in research leading to the development of early RF safety guidelines.

359 Over the past 20 years, dosimetric techniques have been transformed by the use of voxel-based
360 anatomical human-body models and the application of computational methods for the calculation of internal
361 electromagnetic quantities. A voxel (volume pixel) represents a small volume element (or cube) of a tissue with
362 dimensions as small as a few millimetres on each side. Thus, a whole-body human voxel model can consist of
363 many million voxels. The computational algorithms are usually based on the Finite Difference Time Domain
364 (FDTD) method, with some researchers writing their own codes and others using commercially available codes.

365 Several laboratories have developed whole-body animal and human voxel models (Dawson, Caputa &
366 Stuchly, 1997; Dimbylow, 2005a; b; Dimbylow, 1997; Furse & Gandhi, 1998; Lee et al., 2006; Mason et al.,

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367 2000b; Nagaoka et al., 2004). The most widely used voxel model in RF dosimetry was developed by the Brooks
368 AFB Laboratory and based on the database of the Visible Human Project (VHP). Mason et al. (2000a)
369 investigated the effect of different dosimetric parameters on the characteristics of absorption in the VHP Man.

370 Other whole body human voxel models have been developed that are more representative of the
371 average height and weight values of people, as specified by the International Commission on Radiological
372 Protection (ICRP, 2002) and other organisations. Female phantoms have also been developed, e.g. by Dimbylow
373 (2005b) and by Nagaoka et al. (2004).

374 Dosimetry studies have now become more advanced and inclusive, with the development of whole-
375 body human voxel models in different postures, as well as those of children, and pregnant women containing
376 fetuses and embryos. Some of these models were developed by deforming standing adult human models (Cech,
377 Leitgeb & Padiaditis, 2007; Dimbylow, 2006; Findlay & Dimbylow, 2005; Kainz et al., 2003; Nagaoka et al.,
378 2007; Wang et al., 2006). However, whole-body child models have now been developed based on datasets from
379 actual children (Christ et al., 2010b; Lee et al., 2006).

380 Realistic numerical models are developed from images taken by MRI or CT scans. MRI images have
381 generally better quality than those obtained from CT scans, and are useful in identifying interior tissues because
382 high contrast images of different soft tissues can be formed. In order to develop a voxel model for FDTD
383 calculation, original gray-scale images from the scans are interpreted into tissue types, referred to as
384 segmentation. Since the gray scales in MR images do not correspond to tissue types directly, the tissue and organ
385 identification processing has to be performed manually to a large extent. Even if software for automatic
386 identification is applied, manual verification or correction is required. Recent models of the whole human body
387 have up to about 50 tissue types, and the finest resolution is about 1 mm.

388 Intrinsic limitations of voxel-based models are that they do not contain any spatial information at
389 scales smaller than their native resolution and that they cannot be easily deformed to adopt different postures.
390 These problems can be overcome by adopting a computer aided design (CAD) approach to develop models in
391 which the organ and tissue boundaries are represented by parametric surfaces. Christ et al. (2010b) developed a
392 CAD-based human model which can easily move and rotate in any direction with 3-D CAD software and no
393 limitation of their spatial resolution. The surfaces of the model can be readily deformed, but care must be taken
394 for the joints of the body to be correctly articulated. CAD models are usually segmented into voxel models at the
395 required resolution before carrying out FDTD calculations.

396 It is notable that, although many sophisticated and complex numerical models have been developed in
397 recent years, most results of the realistic voxel models have generally agreed with those of the simple anatomical
398 models of the whole body in terms of the whole-body SAR that results from a given plane-wave exposure
399 condition.

400 **3.3.3 Dielectric properties of tissues**

401 Biological tissues contain free and bound charges, including ions, polar molecules and an internal
402 cellular structure. When an external electric field is applied to tissues, the electric charges are shifted from their
403 original position, causing polarisation and ionic drift, and the establishment of displacement and conduction
404 currents. Dielectric properties of tissues (permittivity and conductivity) are measures of these effects which
405 determine the interaction of electric fields with human body. The availability of dielectric data for different
406 tissues is therefore vital for both experimental and computational dosimetry techniques.

407 Dielectric properties of tissues are frequency-dependent and arise as a consequence of particular
408 mechanisms of polarisation, each of which has a maximum frequency at which it can occur. The effect of each
409 mechanism is to cause a “dispersion” in the frequency spectrum of the tissues electrical parameters (the
410 dielectric spectrum). Biological tissues are affected by three main dispersions predicted by known interaction
411 mechanisms. At low frequencies the α dispersion is associated with ionic diffusion processes at the site of the
412 cellular membrane. At intermediate frequencies (kHz region), the polarisation of the cellular membrane, which is
413 a barrier to the flow of ions between the intra and extra cellular media, is the main cause of the β dispersion.
414 Finally, the γ dispersion around the GHz frequency regions mainly due to the polarisation of water molecules
415 inside the tissues. Dielectric properties of tissues vary according to the physiological state of the tissue, the
416 intactness of the cellular membrane and the water content. The α and β dispersions would have importance in the
417 present RF context only if demodulated low frequency signals could arise in tissues through a non-linear

418 mechanism, and if these low frequency signals were strong enough to act through a mechanism with a biological
419 outcome (see Chapter 4).

420 Several techniques exist to assess the dielectric properties of biological tissues. One widely used
421 procedure involves the use of coaxial contact probes and measurement of the reflection coefficient followed by
422 numerical deduction of the dielectric properties of the sample based on a theoretical model of the impedance of
423 the probe (Gabriel, 2000).

424 The early literature on the dielectric properties of body tissues has been reviewed by Gabriel et al.
425 (Gabriel, Gabriel & Corthout, 1996) who also produced their own data (Gabriel, Lau & Gabriel, 1996a) and
426 developed parametric models to reproduce the relative permittivity and conductivity as a function of frequency
427 (Gabriel, Lau & Gabriel, 1996b). This database has been used extensively in dosimetry studies since it was
428 created. Recently, new dielectric data have been collected from live porcine tissues, which are thought to be a
429 good animal-based substitute for human tissue. This also provides an added dimension to the comparison with
430 the data from the 1996 database that were mostly derived from measurement on excised ovine tissue (Gabriel,
431 2005; Peyman et al., 2007).

432 Until recently, the literature data consisted mostly of dielectric properties of tissues from mature
433 animals with a few older studies reporting systematic changes in the dielectric properties of ageing brain tissues
434 (Thurai et al., 1984; 1985). More recent studies looked at the variation of dielectric properties of several rodent,
435 porcine and bovine tissues as a function of animal age and found similar trends (Gabriel, 2005; Peyman,
436 Rezazadeh & Gabriel, 2001; Peyman et al., 2007; Peyman et al., 2009; Schmid & Überbacher, 2005). The results
437 of these studies generally showed a significant decline with age in both permittivity and conductivity of tissues
438 with high water content such as long bone, skull, skin, muscle and bone marrow. At around the GHz frequency
439 region, the observed variations are mainly due to the reduction in water content of tissues as animals age. In the
440 case of brain, increased myelination and decreased water content as a function of age is suggested to be the
441 reason for the drop in permittivity and conductivity values of white matter and spinal cord as animals age, as
442 observed by Schmid and Überbacher (2005) and Peyman et al. (2007). The largest variation in the dielectric
443 properties as a function of age is observed in bone marrow tissues due to the transformation of high water
444 content red marrow to high fat content yellow marrow as the animal grows (Peyman et al., 2009).

445 These studies raised a question over the extent to which the variation of dielectric data as a function of
446 age may affect the results of dosimetry in animal exposure studies, and consequently, the possible implications
447 for studies assessing the exposure of children. The dielectric spectra of ageing porcine tissues have been
448 parameterized to allow their use at any frequency or multiple frequencies within the confines of the models of
449 children (Peyman & Gabriel, 2010).

450 **3.4 Parameters affecting SAR**

451 The effect on SAR of frequency, dielectric parameters, voxel size, polarisation of the incident electric
452 field, shape and size of the model, and posture is reported in several papers, for example, Mason et al. (2000a);
453 Findlay and Dimbylow (2005) and Uusitupa et al. (2010), and have been mentioned in Section 3.1.4.1.

454 Until recently, the majority of dosimetric studies were focused on models of human adults. However,
455 increasing use of telecommunication devices amongst the young has led to research into the differences in the
456 exposure of children and adults.

457 The two main factors differentially affecting the exposure of children and adults are changes in the
458 physical size of the body and in the dielectric properties of tissues (see above). The earlier dosimetric studies
459 used adult human models and scaled them down linearly to approximate a child of a given height and mass while
460 using the dielectric properties of adult tissues (Gandhi, Lazzi & Furse, 1996; Gandhi & Kang, 2002; Wang &
461 Fujiwara, 2003). Later studies developed children's head models from magnetic resonance images of children to
462 obtain a more realistic assessment of the exposure of children (Anderson, 2003; Christ & Kuster, 2005;
463 Dimbylow & Bolch, 2007; Lee et al., 2009; Schönborn, Burkhardt & Kuster, 1998).

464 In considering the dosimetry results from models that use linearly scaled adults to represent children,
465 it is important to recognise that the dimensions of different body structures grow at different rates during the
466 course of development. For example, children have proportionately larger heads than adults. Thus, overall body
467 morphology changes with age, and dosimetry results from scaled adult models should not generally be
468 considered reliable for children, irrespective of dielectric considerations.

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469 **3.4.1 Whole body and localised SAR**

470 As explained above, the thermal effect is the dominant established biophysical mechanism in leading
471 to biological and health effects of RF exposures, and the current guidelines on limiting human exposure are
472 based on limiting heating. Elevation of deep body temperature is closely related to the energy absorption rate in
473 the whole body, or whole-body average SAR (ICNIRP, 1998). Thus dosimetry of RF exposure is generally
474 equivalent to the determination of SAR in the body exposed to RF fields.

475 In the case of extremely localised exposures on some body part, a biologically significant temperature
476 rise could occur around the exposed part resulting in thermal injury of the tissue regardless of the deep-body
477 temperature elevation. Local SAR in the part of the body should be considered in this case. Temperature
478 elevation in the body part, however, is not necessarily proportional to the local SAR because of the heat
479 conduction. Thus dosimetry of RF exposure sometimes includes measurement or estimation of temperature as an
480 adjunctive dose metric since it is more directly related to thermal injury.

481 **3.4.2 Age and size related variations**

482 The effect of age on the calculated SAR could be due to the difference in dielectric properties of
483 younger and older tissues as well as changes in the size and shape of the body/head.

484 **3.4.2.1 Variation in dielectric properties**

485 Some studies have used measured dielectric properties of tissues as a function of age to calculate SAR
486 values in models of children's head/body as a result of exposure to EMFs (Christ et al., 2010a; Gabriel, 2005;
487 Peyman et al., 2009). A few other studies have used adult dielectric properties and adjusted them for younger
488 tissues, assuming higher water content (Dimbylow, Bolch & Lee, 2010; Keshvari, Keshvari & Lang, 2006;
489 Wang et al., 2006). Each of the following three studies have used a different metric in their SAR calculation,
490 such as localised SAR_{10g} values for individual tissues, SAR averaged over 10 g of tissues in the shape of a cube
491 and whole body SAR, therefore a direct comparison cannot be drawn.

492 Gabriel (2005) used age-related dielectric properties of rat tissues in a numerical study of exposure of
493 rat models to plane waves at 27, 160, 400, 900, and 2000 MHz. A total of 34 tissue types have been examined in
494 three rat models (10, 30 and 70 days old), from which only 9 tissues exhibited variation in their dielectric
495 properties as a function of age. The results showed that, although changing the tissue dielectric properties would
496 affect the localised SAR, no clear pattern could be established. The effect on whole body SAR was reported to
497 be small, its extent depending on the variation in properties and the abundance of the tissues in the exposed
498 model. These results can be explained as due to the fact that changes in dielectric properties would affect the
499 coupling with the body and the interaction of tissues with the EMFs. It is also important to isolate the effect of
500 changing tissue properties from all other factors that would affect the exposure, such as the size of the animal,
501 and polarization and direction of the incident field (Gabriel, 2005).

502 The second study examined the sensitivity of calculated SAR averaged over a 10 g cube in one adult
503 and two child head models exposed to EMF from walkie-talkie devices operating at 446 MHz (Peyman et al.,
504 2009). Head models representing adults and 3- and 7-year old children were considered with tissue dielectric
505 properties taken from 10 kg, 50 kg and 250 kg pigs, which were taken as representative of children in the range
506 1-4 years old, children in the range 11-13 years old and adults respectively. The results showed that variations in
507 SAR_{10g} were less than 10% for the investigated configuration and that the variations of the tissue properties are
508 not clearly reflected as a variation in SAR_{10g}. This could be because spatial averaging of the SAR dilutes the
509 effect of changes occurring in masses smaller than 10 g. In addition, different head tissues do not contribute
510 equally in the cubical averaging volume and not all tissues in the averaging volume have the same variation of
511 the dielectric properties with age (in this case only skin contributed to the variation within the 10 g cube).

512 Finally, Christ et al. (2010a) studied the exposure of three anatomical head models (adult, 3- and 7-
513 year old) to a generic dual band mobile phone operating at 900 MHz and 1800 MHz. They incorporated 16 tissue
514 types in the models by assigning dielectric properties of 10 kg, 50 kg and 250 kg pigs. Although the results
515 showed SAR variations due to the age-dependent changes to be within 30%, for all the configurations analysed,
516 age dependencies of dielectric tissue properties did not lead to systematic changes of the 10 g peak spatial SAR.
517 In other words, the hypothesis that variation in the dielectric parameters results in larger exposure of young
518 mobile phone users could not be confirmed. The authors suggested that this may be due to the fact that highest

519 age-dependent variations occur in tissues with low water content, where the SAR would be lower than in the
520 higher water content tissues.

521 The dielectric properties of bone marrow exhibit the largest variation with animal age, as reported by
522 Peyman et al. (2009). As the distribution of bone marrow in the skull is very complex, Christ et al. (2010a)
523 examined the impact of age dependent changes of dielectric properties with the help of a generic planar layered
524 model consisting of skin, fat, bone, marrow, bone again and then muscle. The authors calculated the 10 g peak
525 spatial SAR averaged over bone marrow tissue. The absorption in the bone marrow was largely independent of
526 the layer thickness and clearly age dependent. The results showed that exposure of the bone marrow of children
527 might exceed that of adults by about a factor of ten. This is due to the strong decrease in electric conductivity of
528 this tissue with age. These results are yet to be confirmed in anatomically realistic models exposed to realistic
529 sources.

530 3.4.2.2 Variation in shape and size of head

531 When a small transmitter like a mobile phone is held to the head, there are various general
532 considerations that relate to the distribution of SAR that results in the head. The SAR distribution will be
533 maximal in proximity to the radiating part of the transmitter and will generally reduce through the head as
534 distance increases. Thus, the SAR averaged over the entire head (or brain) will be higher for smaller heads than
535 with larger heads exposed under the same conditions. The radius of curvature of the head in the vicinity of the
536 transmitter is important because a small radius of curvature, as with the child's head, will mean less of the head
537 tissue will be in the immediate proximity of the transmitter and tend to produce lower averaged SAR in the head
538 as a whole. However, the ear acts as a spacer, increasing distance of the transmitter from the head and therefore
539 reducing averaged SAR in the head. Thus, a thinner and more compressible ear, as with a child, will result in
540 higher averaged SAR in the head.

541 The size of the radiating part of a transmitter (antenna) is also a factor affecting the SAR distribution
542 in the head because a larger source will produce a more diffuse input of RF energy to the head tissues than will a
543 smaller source. Overall, antennas reduce in size with increasing frequency and mobile phones have also reduced
544 in size as the technology has developed. Both of these considerations could lead to hypotheses about how
545 localised SAR averaged over 1 g or 10 g might vary with frequency and how they might have changed over time.
546 However, in terms of age-related variations, the situation is complicated, with various factors that can trade-off
547 against each other such that it is difficult to determine without investigation whether localised 1 g or 10 g SAR in
548 the head from a mobile phone would typically be greater or smaller in a child's head than in that of an adult.

549 Gandhi et al. (1996) were amongst the first to use a linearly-scaled child head model with adult
550 dielectric properties to compare the exposure of adults and children to mobile phones. They found a deeper
551 penetration and increase of 50% in the maximum 1 g averaged SAR in the child's head.

552 Hombach et al. (1996) examined the energy absorption in a human head at 900 MHz, and reported
553 that the spatial peak SAR is affected by the size and the shape of human head at a defined distance from source.
554 On the other hand, the authors reported that the effects caused by the complex anatomy are minor in the case of
555 volume-averaged values because these do not exceed the results found in a homogeneous head model.

556 Schönborn et al. (1998) used a scaled adult model and two realistic child head models using MRI data
557 for similar comparisons, but did not report any significant differences between the exposure in adults and
558 children. Bit-Babik et al. (2005) reported that the peak local average SAR over 1 g and 10 g of tissue and the
559 electromagnetic penetration depths are about the same in head models of adults and children under the same
560 exposure condition.

561 De Salles et al. (2006) reported that under similar exposure conditions (850 and 1850 MHz), the 1 g
562 SAR calculated for children is higher than that of adults. The authors claimed that when using a 10-year old
563 child model, SAR values are 60% higher than those obtained for adults.

564 Attempts to repeat the above studies by other researchers have produced contradictory results, in
565 particular when different normalisations were performed for output power or antenna current. In addition,
566 differences in averaging procedures used to determine spatial peak SAR meant there was a need for further
567 studies using a standardised approach (Bit-Babik et al., 2005; Wang & Fujiwara, 2003). Some of the subsequent
568 studies reported a higher SAR in the head models of children (Anderson, 2003; Wiart et al., 2008) while others
569 did not find any significant difference in SAR between adult and child models or could not reach a conclusion

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570 (Christ & Kuster, 2005; Fujimoto et al., 2006; Keshvari & Lang, 2005; Lee, Choi & Choi, 2007; Martínez-
571 Búrdalo et al., 2004; Wang et al., 2006).

572 In 2006, a large study was carried out involving 14 laboratories to study the differences in SAR for
573 adults and children who are exposed to mobile phones (Beard et al., 2006). A single protocol was followed by all
574 the participants using the Specific Anthropomorphic Mannequin (SAM) homogeneous head model designed for
575 mobile phone compliance measurements, an anatomically realistic adult head model and a scaled 7-year old
576 child head model. Two operating frequencies of 835 and 1900 MHz and two phone positions, denoted cheek and
577 tilt, were considered. In addition, the SAR values were normalised to both the antenna input power and the feed
578 point current. The results revealed a reverse effect at 1900 MHz, where the peak 1 g SAR values in the head of
579 the adult model were higher in both head positions and for both normalisation scenarios. At 835 MHz however,
580 the SARs were higher in the child model than the adult model in particular in the tilt position and when
581 normalizing to antenna current.

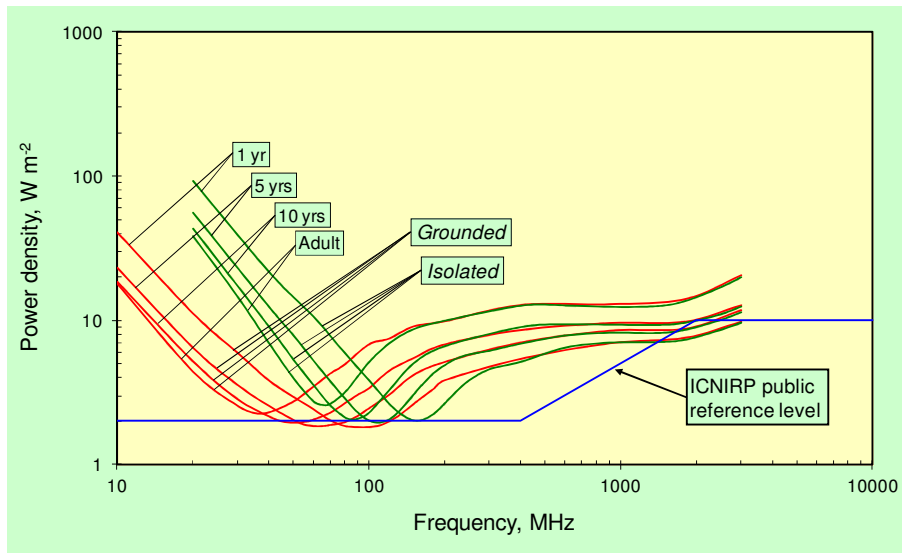
582 Hadjem et al. (2005) found that brain exposure depends on the morphology of the ear and Wiart et al.
583 (2008) reported higher exposure of the cerebral cortex in children than in adults.

584 In summary, the overall evidence from studies on the exposure of children's heads does not draw a
585 consistent picture to support the assumption that the SAR level in a child's head is higher than that for adults.
586 This arises because the results are highly model-specific and power absorption in the head is influenced by many
587 factors. It is also notable that even the higher SAR values reported in some of the studies for children's heads,
588 are below the 2 W/kg ICNIRP basic restriction.

589 3.4.2.3 *Variation in shape and size of body*

590 The effects of body size in relation to the wavelength have been discussed in general terms earlier in
591 this chapter, and pertinent studies are reviewed here. As with the studies above regarding the head, the
592 differences between the exposure of children and adults are an important theme in the literature.

593 It is established that the whole body SAR depends on the size of the body, and this has been
594 considered when the ICNIRP reference levels were set. Starting from the premise that an SAR of 4 W/kg for a
595 healthy adult is equivalent to a 1°C temperature rise, there is a 50-fold reduction in deriving the basic restriction
596 on whole body SAR for the general public (0.08 W/kg). In principle, this reduction factor should be sufficient to
597 take into account all the variations due to different dosimetric factors, including body size. However, it is now
598 apparent that the basic restriction on exposure for small children under worst case exposure conditions may be
599 slightly exceeded at the public reference level (NRPB, 2004a; 2004b). This is supported by studies carried out
600 for phantoms modelling children in the age range from 9 months to 11 years, which found that whole body SAR
601 restrictions are exceeded in two frequency ranges: 45–170 MHz and 1400–4000 MHz (Bernardi et al., 2003;
602 Dimbylow & Bolch, 2007; Dimbylow, 1997; 2002; Mason et al., 2000a; Tinniswood, Furse & Gandhi, 1998;
603 Wang et al., 2006). Dimbylow (2007) reported that, when using a phantom of a 9-month old baby at a frequency
604 of 1.6 GHz, the whole body SAR was equal to the basic restriction when the electric field strength was 0.83
605 times the value of the reference level. Thus, exposure at the reference level would result in the whole body SAR
606 restriction being exceeded by about 40%, given that SAR is proportional to the electric field strength squared.
607 Higher SAR levels were also observed in phantoms modelling children aged 3, 5, and 7 years in the frequency
608 region around 2 GHz (Conil et al., 2008; Hirata, Asano & Fujiwara, 2007; Neubauer et al., 2009). Figure 3.3
609 shows a typical set of results illustrating that the reference level is not conservative over the basic restriction in
610 certain frequency ranges with smaller body sizes.



611

612 **Figure 3.3.** Curves showing the power density of a vertically polarized plane wave incident on the front of the
 613 body required to produce a whole-body SAR equal to the basic restriction with different size phantoms and
 614 different grounding conditions (Dimbylow, 2005a; 1997).

615 The polarization of the incident wave can also affect the calculated SAR, as studied by Hirata and
 616 Fujiwara (2009) on whole body averaged SAR in an infant model. The authors reported that whole body
 617 averaged SAR for plane-wave exposure with vertically aligned electric field is smaller than that with a
 618 horizontally aligned field for frequencies above 2 GHz. This could be due to the component of the surface area
 619 perpendicular to the electric field of the incident wave (Hirata & Fujiwara, 2009).

620 Uusitupa et al. (2010) also suggested that the effect of polarization must not be neglected. They
 621 reported that near the whole-body resonance, polarization has a strong effect on whole body SAR. They also
 622 indicated that in the GHz region the effect may be more than 2 dB (i.e. a variation of 60%). Note, between 2 and
 623 5 GHz for adults, whole body SAR is higher for horizontally polarised fields than for vertically polarized ones if
 624 the incoming direction is in the azimuth plane. For a child, however, the effect of incoming direction is similar to
 625 that of an adult, except at 300 MHz for horizontal polarization. Generally, it is suggested that in the GHz range
 626 horizontal polarization gives higher whole body SAR than with vertically aligned fields (Kühn et al., 2009;
 627 Vermeeren et al., 2008).

628 In summary, most of the studies conducted on whole body SAR, reported higher exposure for child
 629 models than for adults in two frequency regions (around 100 MHz and 1–4 GHz). They also show that, within
 630 these specific frequency ranges, an exposure to an electric field strength or plane wave equivalent power density
 631 equal to the reference level can lead to whole body SAR exceeding the basic restriction by up to 40% under
 632 worst case conditions. Thus, for children or small persons (shorter than 1.3 m) the reference levels of the
 633 ICNIRP recommendations may not be conservative estimates. A review by ICNIRP (2009) acknowledges the
 634 outcome of the above studies and concludes that this increase in SAR is negligible compared with the large
 635 reduction factor of 50 applied in deriving the whole-body basic restriction for the general public (a reduction
 636 factor of around 30 still remains). Similar conclusions were reached in comprehensive reviews carried out by
 637 national authorities such as the Health Council of the Netherlands (HCN, 2011).

638 **3.4.2.4 SAR in the fetus**

639 Dimbylow (2007) reported that the whole body SAR is lower in pregnant women compared with non-
 640 pregnant women and that the difference increases with gestation period. This is because the whole body
 641 averaged SAR depends inversely on the mass, which increases during pregnancy while the height remains
 642 constant. The study also found that the fetus's maximum whole body SAR occurs at 70 MHz for conditions
 643 where the mother is isolated from ground and SAR levels in the fetus are lower compared with those in the
 644 mother. The results also confirmed that the ICNIRP reference levels are sufficient to guarantee compliance with
 645 the basic restriction for plane wave exposure. Two more studies reported similar results and confirmed
 646 compliance with ICNIRP guidelines (Nagaoka et al., 2007; Togashi et al., 2008). Kawai et al. (2006) developed
 647 a model of the abdomen of a pregnant woman based on MR images of Japanese women in late pregnancy period.

648 They then assessed the local 10 g averaged SAR of a foetus inside this abdomen to be less than 1.5 W/kg when
649 exposed to 150 MHz signals with the maximum power of 5 W used for portable radio terminals in Japan. The
650 same authors (Kawai et al., 2010) reported the maximum average SAR in the embryos exposed to plane waves to
651 be lower than 0.08 W/kg when the incident power density is equal to the ICNIRP public reference level.

652 **3.4.3 Posture and grounding effects**

653 A change in the posture of the human body can significantly affect the way in which it absorbs RF
654 EMFs. Findlay and Dimbylow (2005) studied the effects of posture on FDTD calculations of specific absorption
655 rate in an anatomically realistic model of the body which was modified to develop new voxel models in postures
656 other than the most commonly analysed standing position with arms to the side. The scenarios considered were
657 sitting, arms stretched out horizontally to the sides and arms raised vertically above the head. The results showed
658 that the effect of a raised arms posture lowered the resonant frequency and increased the value of the whole-body
659 averaged SAR at resonance by up to 35% when compared to the reference arms by the side position. They also
660 reported that in certain postures external electric field reference levels alone would not provide a conservative
661 estimate of localised SAR exposure and it would be necessary to invoke secondary reference levels on limb
662 currents to ensure compliance with ICNIRP's basic restrictions. A similar study was carried out for SAR in a 7-
663 year old child voxel model (Findlay, Lee & Dimbylow, 2009). It reported that raising the arms increased the
664 SAR by 25%.

665 Uusitupa et al. (Uusitupa et al., 2010) suggested that body posture has little effect on whole-body SAR
666 in the GHz region, but at around 300-450 MHz, one may even expect a 2 dB rise (+60%) in whole-body SAR if
667 posture is changed from the default standing position to a sitting position. The authors also concluded that
668 posture affects peak 10 g SAR (in head and trunk and limbs) much more than whole-body SAR.

669 **3.5 Temperature elevation**

670 Temperature elevation can be considered as the main factor able to cause adverse health effects in the
671 human body when it is exposed to external RF EMF. It results from energy deposited or absorbed in body tissues
672 through local, partial-body or whole-body exposures.

673 The extent of tissue temperature rise depends on several factors including the pattern of energy
674 absorption, the pathways through which heat is transferred and removed from the tissues inside the body, heat
675 exchange between the body surface (skin) and the external environment, and the thermoregulatory process.

676 It is impossible to non-invasively measure the temperature rise inside the body; however
677 computational simulations based on the internal distribution of SAR can be used to make estimations.

678 Pennes (1948) proposed the so called "bioheat" equation to calculate the time variation of
679 temperatures in a human body. It gives the temperature distribution, $T(r,t)$, which is a function of location and
680 time inside the body taking into account the SAR distribution, $Q_v(r)$, expressed as SAR divided by the volume
681 density $\rho(r)$.

$$682 \quad \nabla \cdot (K(r)\nabla T) + A(r,T) + Q_v(r) - RL(r) - B(r,T)(T - T_B) = C(r)\rho(r) \frac{\partial T}{\partial t} \quad [\text{W/m}^3] \quad (3.3)$$

683 The bioheat equation takes into account several factors such as the tissue thermal conductivity,
684 K (W/m °C), metabolic heat production A (W/m³), respiratory heat losses from lungs RL (W/m³), difference
685 between blood and tissue temperature ($T - T_B$) and the heat exchange due to capillary blood perfusion
686 B (W/°C per m³), which is proportional to blood flow.

687 Originally, the bioheat equation did not account for thermoregulatory response, and thermal responses
688 were later considered by using very simplified models of human bodies (Stolwijk & Hardy, 1977). Foster and
689 Adair (2004) tested the response of the model by using experimental data from human volunteers. Finally,
690 Bernardi et al. (2003) incorporated the thermal response model within the bioheat equation; which can be used to
691 estimate the temperature elevation in an anatomically-based human body model in the time domain.

692 Although the bioheat equation estimates the temperature distribution in the body as a function of time,
693 for pulsed or brief applications of RF energy, the exposure duration is not long enough for significant conductive

694 or convective heat transfer to affect tissue temperature rise. In this case, the initial rate of rise in temperature is
695 related to SAR as follows:

$$696 \quad \Delta T = \frac{SAR}{C} \Delta t \quad (3.4)$$

697 where ΔT is the temperature increment ($^{\circ}\text{C}$), C is the specific heat capacity of tissue ($\text{J}/\text{kg}^{\circ}\text{C}$), and Δt
698 is the duration of RF exposure. Equation (3.4) shows that the rise in tissue temperature during the initial
699 transient period of RF energy absorption is linearly proportional to SAR and inversely proportional to the
700 specific heat capacity of tissue.

701 Clearly, if the intensity and duration of the exposure are excessive, the RF energy can produce
702 temperature rises that can result in adverse health effects due to heat stress or local tissue damage.

703 Under moderate conditions, a temperature rise on the order of 1°C in humans and laboratory animals
704 can result from an SAR input of approximately $4 \text{ W}/\text{kg}$. However, such temperature rise falls within the normal
705 range of human thermoregulatory capacity. Above this temperature or SAR value, disruption of work in trained
706 rodents and primates has been reported for normal environmental conditions (ICNIRP, 1998).

707 **3.5.1 Localised exposure and thermal time constants**

708 In estimating the localised temperature elevation in the human body as a result of localised exposure,
709 one needs to consider the balance between the rate of RF power deposition and the time constants for heat
710 convection by blood flow and heat conduction. For sufficiently high whole-body and intense-localised exposure,
711 the absorbed EM energy becomes appreciable in comparison with basal metabolism (around 1 to $2 \text{ W}/\text{kg}$
712 typically) and leads to body-core temperature elevation (Adair, Mylacraine & Allen, 2003; Guy et al., 1975). A
713 proportion of the RF energy absorbed locally is transferred as heat to the body core via blood flow, which leads
714 to the activation of centrally mediated (e.g. through sensors in the brain) thermoregulatory responses in order to
715 maintain core body temperature (Adair & Black, 2003).

716 The time taken for blood to be travel from a site of localised exposure to where its temperature is
717 sensed centrally means that the thermal time constant of the body core is slightly longer than that associated with
718 temperature elevation in a body part due to localised exposure. Moreover, thermal sensors are also located in the
719 skin, which means that an adaptive thermoregulatory response often begins before an increase in core
720 temperature. This is particularly the case when RF energy absorption is mostly confined to the surface of the
721 body, as occurs with exposures towards the higher RF frequencies.

722 In addition, certain tissues, e.g. the brain, have higher rates of blood flow than other tissues, even
723 when body core temperature is not significantly increased, and this mechanism can play an important role in
724 limiting temperature rises for intense localised exposure (Wainwright, 2003).

725 **3.5.2 Whole body exposure and thermal time constants**

726 When estimating the whole body temperature elevation, it is important to take into account the
727 variation in blood temperature. Bernardi et al. (2003) incorporated the blood temperature factor into the original
728 bioheat equation. For example in the case of a plane-wave exposure at 40 MHz with a power density of
729 $2 \text{ W}/\text{m}^2$ (equal to the ICNIRP reference level) the maximum temperature rise at the ankle is about 0.7°C .

730 As mentioned before, the body's thermoregulatory response reduces temperature rise, particularly in
731 the body core. Hirata et al. (2007) considered the far-field (plane wave) exposures of the naked body at two
732 frequencies, 65 MHz (whole body resonance frequency) and 2 GHz , taking into account the effect of
733 perspiration on body-core temperature elevation. The variability of temperature elevation caused by sweating
734 was found to be 30% . A whole-body average SAR of $4.5 \text{ W}/\text{kg}$ was required for a body-core temperature
735 elevation of 1°C after 60 min exposure in the model of human with the lower sweating coefficient, while
736 $8 \text{ W}/\text{kg}$ was required with the highest sweating coefficients. The thermal time constant in the body core was 20
737 min and almost the same for both frequencies.

738 **3.5.3 Temperature rise in the eye**

739 Excessive temperature rise in the eyes is associated with several effects, including cataract formation.
740 Guy et al (1975) reported microwave-induced cataract formation in rabbit eyes. Emery et al. (1975) developed a
741 heat transfer model for the rabbit eye by assuming that the eye is an object thermally isolated from the rest of the
742 head. Later works used anatomically based human models and improved heat transfer coefficients between the
743 eye, the air and the rest of the head to quantify the temperature rise in the eye (Lagendijk, 1982).

744 Two studies reported a maximum temperature rise of 0.04–0.06 °C on the surface of the eye and in the
745 lens when they are exposed to an incident power density of 10 W/m². The first study considered the frequency
746 range 0.6–6 GHz (Hirata, Matsuyama & Shiozawa, 2000) and the second considered 6–30 GHz (Bernardi et al.,
747 1998). These studies also confirmed that the maximum temperature elevation in the lens decreases with
748 increasing frequency.

749 Several other authors who used improved heat transfer models and took into account the blood flow in
750 the choroidal and retinal tissues as well as the heat transfer in the whole head reported a temperature rise of
751 0.3°C for an eye-average SAR of 2 W/kg (Buccella, De Santis & Feliziani, 2007; Hirata, 2005; Wainwright,
752 2007). As expected, a correlation was observed between the average eye SAR and the maximum temperature
753 elevation in the lens. However, a lower temperature elevation was reported in Flyckt et al. (2007) using a heat
754 transfer model involving discrete vasculatures (DIVA). Van Leeuwen et al. (1999) also used DIVA modelling
755 and reported lower local temperature elevation around the blood vessel due to the cooling effect of blood flow.

756 **3.5.4 Temperature rise in the head**

757 Following increasing use of telecommunication devices, a considerable number of studies have been
758 carried out on the temperature elevation in the head exposed to mobile phone handset antennas (Bernardi et al.,
759 2000; 2001; Gandhi, Li & Kang, 2001; Hirata, Morita & Shiozawa, 2003; Hirata & Shiozawa, 2003; Hirata et
760 al., 2006; Ibrahim et al., 2005; van Leeuwen et al., 1999; Wainwright, 2000; Wang & Fujiwara, 1999). Although
761 some common assumptions were made, each of the above studies used a different type of antenna, head model
762 and SAR metric, therefore a direct comparison between the results is not possible. The output power of a typical
763 handset antenna is up to a few hundred milliwatts, which is much lower than the basal metabolic rate of an adult
764 male of 100 W or more; therefore the above studies have all assumed a constant temperature for blood.
765 Wainwright (2000) used a finite element thermal model to allow better simulation of the surface curvatures of
766 the human head to calculate SAR and temperature elevation. Hirata and Shiozawa (2003) studied the correlation
767 between peak spatial-average SAR and maximum temperature elevations in the head for different frequencies,
768 polarizations, feeding positions, and antennas. The authors found fairly good correlations between peak spatial-
769 average SAR and maximum temperature elevation in the head excluding the pinna. Hirata et al. (2006)
770 investigated the correlation of maximum temperature rise in the head with peak SAR calculated by different
771 average schemes and masses. For an exposure scenario involving a peak SAR of 10 W/kg (for 10 g of
772 contiguous tissue and for exposure times of 60 min or longer), including and excluding the pinna, resulted in a
773 maximum temperature rise of 2.4 or 1.4 °C in the head respectively. The authors also noted that a high degree of
774 spatial correlation between peak SAR and maximum steady state temperature rise for durations of 60 min or
775 longer is not expected, especially for exposures of large biological bodies with efficient thermal transfer
776 characteristics (Hirata et al., 2006).

777 Other authors reported a temperature rise of about 2 °C in a head without pinna when peak SAR was
778 10 W/kg averaged over a 10 g cubic volume, which was higher than that for contiguous tissue volumes (Bernardi
779 et al., 2000; Hirata & Shiozawa, 2003; Hirata, Shirai & Fujiwara, 2008; Razmadze et al., 2009; Wainwright,
780 2000).

781 The mobile phone handset (and its battery) placed close to the head can also behave as a source of
782 conducted heat and reducer of convection from the skin surface, thereby causing local temperature rise (Bernardi
783 et al., 2001; Gandhi, Li & Kang, 2001; Ibrahim et al., 2005). The skin temperature elevations due to these
784 processes have been shown to be 1°C or more, which is greater than that caused by RF energy deposition.

785 **3.6 Contact and induced currents**

786 At frequencies lower than 10 MHz, where the wavelength of RF radiation is at least an order of
787 magnitude longer than the dimensions of the human body, sources giving rise to appreciable exposures tend to
788 be those that are in close proximity to the body and therefore at distances well within a wavelength. In such

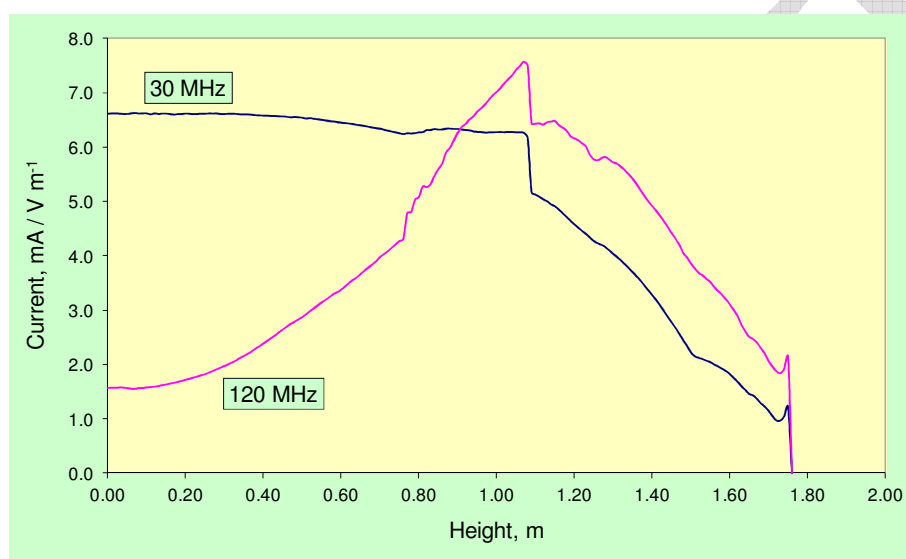
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789 circumstances, coupling and field behaviour inside the body are characterized by near-zone reactive fields and
790 are quasi-static in character. The electric and magnetic fields become decoupled, and they act separately and
791 additively inside tissue medium (Lin, 2000; 2007). The total induced fields can therefore be obtained by
792 combining the two independent quasi-static electric and magnetic solutions of the electromagnetic field theory.

793 For an induced electric field E (V/m) at a point in the body, the corresponding induced current
794 density, J (A/m²), can be calculated through the following formula:

795
$$J = \sigma E \tag{3.5}$$

796 where σ is the electrical conductivity (S/m) of the tissue. It is then possible to integrate J over the body cross-
797 section in vertically stacked horizontal slices in order to gain an understanding of how induced current varies
798 with height through the body. Figure 3.4 shows an example calculation from Dimbylow (1997) at a frequency
799 just below the body resonance (30 MHz) where the maximum layer current occurs at the feet and a frequency
800 above resonance (120 MHz) where the maximum layer current occurs at the mid-height of the body. Note that
801 the step at 1.1 m height arises because the current in the arms has been neglected in the calculation.



802
803 **Figure 3.4.** Variation of layer-averaged current with height in the body for a vertically polarized uniform field
804 incident towards the front of the body.

805 The magnitude of induced currents depends on similar factors to SAR, such as the electric and
806 magnetic field strength, the polarization of the field, and the grounding conditions. In general, the current must
807 reduce to zero at the top of the head, but it will not reduce to zero at the feet, especially if there is a good contact
808 to the ground. Similarly, current may flow through and out of an arm if the hand is in electrical contact with a
809 conducting object.

810 Operators of RF plastic sealers represent an occupational category that is highly exposed to RF EMFs
811 (Wilén et al., 2004), where the induced current flowing from the feet to the ground may reach few hundreds of
812 mA. The coupling of the body to the electric field component from these devices is usually stronger than that of
813 the magnetic field. The existence of high electric fields around the electrodes would induce RF currents that
814 would flow along the legs and torso. Maximum absorption occurs in the limbs (70% of the total absorbed power)
815 where the current density, and therefore also the localised SAR, increases considerably due to a small cross
816 section and a high amount of low-conductivity bone. Computational studies suggest that the local 10 g average
817 SAR is about 10 W/kg for the arm and 5–8 W/kg for the foot with 100 mA current through that limb (Dimbylow,
818 2001; Findlay & Dimbylow, 2005). In the case of good contact with the ground the current in the lower legs,
819 increases considerably and the maximum current density shifts to the ankles. Additionally, the whole body
820 average SAR may increase by a factor of 2 as Chen, Gandhi and Conover (1991) have reported. Gandhi et al.
821 (1997) suggested that other posture-related changes such as extending the hands over the electrodes of sitting
822 may further increase the SAR by a factor of 2 or more (Gandhi et al., 1997). According to Jokela and Puranen
823 (1999), the induced current is not affected strongly by the variation in the electric field as a function of the
824 distance because the whole body integrates capacitive displacement current. As the distance increases, the

825 electric field becomes more uniform which partly compensates for the decrease of the peak field at the electrode
826 plane.

827 **3.7 Auditory effect**

828 Exposure to a high peak power pulsed RF signal may result in a physiologically insignificant but rapid
829 temperature rise leading to an auditory stimulation response in animals and humans. The effect occurs at
830 frequencies ranging from hundreds of MHz to tens of GHz and is known as the microwave auditory (or hearing)
831 effect (Lin, 1980; 2007; Lin & Wang, 2007). Transient localised heating during pulses leads to tissue expansion
832 and generates an acoustic wave of pressure that travels by bone conduction to the inner ear where it activates the
833 cochlea receptors the same way as in normal hearing. A single microwave pulse can be perceived as an acoustic
834 click or knocking sound and a train of microwave pulses to the head can be sensed as a buzz or audible tune,
835 with a pitch corresponding to the pulse repetition rate.

836 Experimental and theoretical studies have shown that the microwave auditory phenomenon does not
837 arise from an interaction of microwave pulses directly with the auditory nerves or neurons along the auditory
838 neurophysiological pathways of the central nervous system. Peak power densities of a few kW/m² are required to
839 exceed the threshold acoustic pressure for hearing in humans (20 mPa).

840 The induced sound frequency exhibits an acoustically resonant behaviour where pulses appropriately
841 separated in time are able to reinforce the pressure waves in the head, leading to increased sensitivity. The
842 resonant frequency depends on the size of the head; the smaller the head radius, the higher the frequency. For
843 rat-size heads, the predicted acoustic frequencies are 25 to 35 kHz in the ultrasonic range, which rats can easily
844 hear. For the size of human heads, resonant frequencies of 7 to 15 kHz have been found which are clearly within
845 the audible range of humans.

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